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Can sonication enhance release from liquid-core capsules with a hydrogel membrane?

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ABSTRACT

The objective is to investigate the influence of sonication on the mechanical and release properties of hydrogel capsules. A new fabrication process is developed to fabricate millimetric capsules made of a highly-viscous liquid core protected by a thin hyperelastic alginate membrane. At high intensities and/ or long exposure times, sonication can lead to the capsule rupture, because it induces fatigue in the membrane. Below the breakup threshold, no remnant effect of sonication is, however, measured on the capsule mechanical properties. The release is studied by sonicating capsules filled with blue dextran suspended in an aqueous solution. The mass release that results from sonication is found to be proportional to the sonication duration time and pressure wave amplitude. A possible physical interpretation is that the acoustic streaming flow induced by the ultrasonic wave enhances convection in the vicinity of the capsule membrane and thus mass release. We have finally quantified the passive release subsequent to low-intensity sonications: it is on average identical to the one measured on non-sonicated capsules. Overall the membrane therefore recovers its physical and mechanical properties after sonication. If sonication leads to an increase in porosity of the capsule membrane, the increase is temporary and reverses back at the end of the ultrasonic stimulation.

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39 1. Introduction

Drug release from polymeric delivery systems responsive to 40 41 external stimulations is receiving increasing attention in therapeutic medicine [1]. It is used to remotely control not only the rate of 42 drug delivery to efficiently meet the time-evolving needs of 43 patients, but also the site of delivery to induce a targeted delivery. 44 The release rate mainly depends on the sensitivity of the drug vec-45 tor to the stimulation. Among possible stimulations, one finds tem-46 perature change [2], pH [3], light exposure [4], magnetism [5]. 47 External force stimulation is another method to induce release 48 from drug carriers. Zanina et al. [6] studied the influence of shear 49 50 stress and found that release is only reversible at low shear stress. Lee et al. [7] observed that when compression is applied on drug-51 loaded alginate gels, the gel reacts like a sponge and the drug is 52 squeezed out of the gel. Such experience has never been conducted 53 54 on capsules.

Less attention has been paid on ultrasonic stimulation. Three decades ago, it has been suggested that ultrasounds could increase the degradation and permeability [8] of polymeric matrices and hence the release of embedded drugs. Currently, ultrasonic stimulation of microbubbles is routinely used in the field of med-

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ical imaging. Microbubbles (<2 µm in size) have been found to be natural innocuous contrast agents, when stimulated by ultrasounds [9]. As the bubbles need to be coated by a lipid layer to be stabilized, they can simultaneously serve as drug vectors for targeted delivery. The drug release results from the large oscillations induced by the ultrasonic stimulation in the gas core. The efficiency of encapsulated bubbles as drug carrier is, however, limited owing to the small quantity of active material that can be carried in the shell and their short lifetime [10]. Liquid-filled carriers offer a good alternative, as drugs are typically aqueous solutions. But, if larger quantities of drugs can be encapsulated, the release mechanism induced by sonication needs to be established for each vector type.

A few studies have tested the effect of sonication on drug release from liquid-filled carriers. Schroeder et al. [11] sonicated liposomes, consisting of a lipid bilayer encapsulating a liquid drug. They showed that sonication induces transitory reversible pores on the liposome, which leads to an increase in the encapsulated drug release. They found that sonication at low frequency (20 kHz) is more efficient than at high frequency (>1 MHz) and that the higher the ultrasonic power density, the larger the drug release. Similar results were obtained on Pluronic micelles, which are spherical structures with a lipid monolayer [12].

Capsules are another type of drug carriers, for which the liquid core is protected by a solid membrane with elastic properties. The membrane can be made of various constituents, e.g. reticulated

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86 polymers [13], reticulated proteins [14], polyelectrolytes [15]. 87 Their self-release ability is determined essentially by the mem-88 brane properties. The use of ultrasonic stimulation to promote 89 the release rate has only been studied on submicron-sized capsules 90 with a rigid polyelectrolyte shell [16]. Drug release only occurs if 91 the particle rigid wall is broken up. Shchukin et al. [16] showed 92 that a low-frequency stimulation (20 kHz) can lead to the capsule 93 breakup and that, for a constant intensity of sonication, the dura-94 tion time of sonication required for breakup increases with the 95 membrane rigidity. But the particles used can be considered rather 96 as liquid-filled rigid microcontainers than as actual capsules, as 97 capsules intrinsically have a deformable membrane by definition.

98 No study has yet considered the effect of sonication on softmembrane liquid-filled capsules. Our present objective is to mea-99 100 sure the effects of sonication on polyelectrolyte capsules in order 101 to understand the release mechanism that can be induced by son-102 ication. As a first exploratory study, we test millimetric alginate 103 capsules. We measure the evolution of the geometrical and 104 mechanical properties varying the sonication parameters. 105 Kühtreiber et al. [17] has indeed shown that the mechanical prop-106 erties of hydrogel capsules play an important role on their release 107 properties and that any damage on the capsule membrane might 108 change the release behavior.

109 The methods used to assess the mechanical properties of cap-110 sule membranes are typically based on the measurement of defor-111 mation under a well-defined stress; a mechanical model of the 112 capsule deformation is needed to infer the membrane elastic prop-113 erties from the experimental data. Large artificial capsule can be 114 subjected to a shear force in a spinning rheometer [18]; the tech-115 nique is, however, limited by the rather low level of mechanical 116 stress that can be applied to the capsules. They can also be 117 squeezed between two rigid parallel plates: both the distance between the plates and the compression force are measured simul-118 119 taneously. This technique is often used to evaluate a bursting force 120 only [19]. It is, however, possible to also extract the membrane 121 mechanical properties through inverse analysis, combining com-122 pression measurements with engineering models [20,21]. The 123 experiments consist in measuring the force needed to subject the 124 capsule to a pre-defined deformation. An analytical or numerical 125 model, based on an assumed constitutive law, is then used to 126 deduce the membrane mechanical properties. In the present study, 127 we will apply this technique following the method developed by 128 Carin et al. [22].

129 The release properties will then be investigated on capsules filled with a 0.5% blue dextran solution. A few methods exist to 130 131 evaluate the release of an encapsulated substance. The fluores-132 cence detection method requires the encapsulated fluid molecules 133 to be labeled by fluorescence. The evolution of the fluorescence 134 intensity is measured either in the encapsulated fluid or in the sur-135 rounding medium in order to estimate the release [23]. High-136 pressure liquid chromatography (HPLC), another technique used 137 to detect chemical substances in a mixture, may also be applied to quantify released substances [24]. But the method we have cho-138 sen is spectrophotometry, which is the most frequently used owing 139 140 to its high precision and simplicity of use [25]. The absorbance value is measured with the spectrophotometer at various instants of 141 142 time in the external solution containing the capsules. The corresponding mass released can be calculated calibrating the measure-143 ments on samples of known concentrations. 144

All the experimental techniques used in the study are described in Section 2. We present the new method developed to fabricate the capsules, as well as those used to expose them to ultrasounds and measure their mechanical properties and release. The results are then presented and discussed in Section 3. We first show the effect of sonication on the capsule mechanical properties and possible breakup, before considering its influence on the release of encapsulated molecules. The release induced by sonication is152compared with the one induced by compression. We finally show153the influence of sonication on the passive release, to study whether154sonication has a permanent effect on the capsule porosity. Conclusions on the influence of sonication on capsule release are provided156in Section 4.157

2. Materials and methods

2.1. Preparation of liquid-core alginate-membrane capsules

A new fabrication method has been designed to produce cal-160 cium alginate capsules. It is inspired from the fabrication processes 161 of Nigam et al. [15] and Nussinovitch et al. [26]: the hydrogel 162 membrane capsules are likewise obtained by extrusion in a one-163 step process. A solution containing 40% (w/v) sucrose of molecular 164 weight 342.3 Da (84100, Sigma Aldrich, USA) and 0.5% (w/v) cal-165 cium chloride serves as the liquid core of the capsule (solution 166 A). Sucrose is used as a non-gelling polymer to constitute the core 167 of the calcium-alginate capsules and ensure their spherical shape. 168 We chose to use sucrose similarly to Nussinovitch et al. [26] and 169 not dextran like Nigam et al. [15], because sucrose has a much 170 smaller molecular weight than dextran; it can therefore be washed 171 off more easily from the capsule core, once the membrane is 172 created. 173

The solution is extruded through a 24 gauge needle (Fisher Scientific) by a peristaltic pump (Ismatec ISM834C, Switzerland) at a flow rate of 1 ml/min. Droplets of solution A form at the tip of the needle and fall into a 0.2% alginate (A0682, Sigma Aldrich, USA) solution (solution B). A distance of 3 cm between the tip of the needle and the surface of solution B ensures spherical droplets. The alginate molecules contained in solution B immediately react with the calcium cations of solution A at the droplet interface leading to the formation of a hydrogel membrane. The reaction time determines the thickness of the membrane. After 5 min, the reaction is stopped by a fivefold dilution of the alginate solution with distilled water. The small size sucrose molecules contained in the capsule core are then cleared off washing the capsules with a large volume of distilled water. The capsules are stored in a 0.5% (w/v) calcium chloride solution (solution C), isotonic with the internal liquid, to stabilize the gel membrane. Tests are conducted after one day of storage.

To fabricate blue dextran filled capsules, blue dextran of molecular weight 2000 kDa (D5751, Sigma Aldrich, USA) is added to solution A at a concentration of 5 mg/ml. The capsule fabrication method remains identical for the rest. The capsules are stored in a modified solution C containing the same concentration in blue dextran as the core solution in order to avoid its diffusion. Before use, they are washed with distilled water to clear off the blue dextran molecules sticking on the capsule surface.

2.2. Capsule dimensions

The capsule dimensions are obtained by capturing images of the
capsules with a CCD camera (JAI M50, Imasys S.A., France). The
images are acquired with the Scion Image software (Scion Image,
Scion Corporation, USA) and analyzed with Image J 1.42q (National
Institutes of Health, USA). The upper piston, 8.0 mm in diameter, is
used as reference length scale for the calibration.200201
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The capsules present a small departure from sphericity and tend to be oblate ellipsoids. We find a 9% difference between their initial height D_0 and width L_0 measured on the images. We have then calculated their volume assuming the capsules to be axisymmetric, and their equivalent radius r_0 , defined as the radius of the sphere

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having the same volume. On average, the capsules have a mean radius $r_0 = 1.35 \pm 0.01$ mm.

212 213 As the capsule membrane appears as more opaque than the li-214 quid core on the images, its thickness can be directly obtained from the acquired images. The thickness is measured at four distinct 215 locations every 90°, in order to take into account possible thickness 216 217 variations around the capsule circumference. The average thickness measured on a set of capsules is found to be $h_0 = 0.20 \pm 0.01$ mm 218 $(h_0/r_0 = 14.8\%).$ 219

2.3. Capsule sonication 220

221 A 30 kHz ultrasonic generator (UP50H, Hielscher, Germany) with a 7 mm sonotrode (MS7, Hielscher, Germany) is used for 222 the capsule sonication. The parameters of sonication that are var-223 224 ied in the study are the duration time t_s and power P_s of the ultra-225 sonic stimulation. The sonication power can be adjusted by changing the oscillation amplitude of the sonotrode. The sonication 226 227 time is ranged from 2 to 35 min, and the sonication power from 228 0.48 to 17.46 W.

229 To determine the influence of sonication on the capsule 230 mechanical properties, three capsules are placed in solution C in 231 a 15 mm diameter tube, itself placed in a water bath to avoid 232 any temperature increase induced by sonication. The sonotrode tip is immersed in the solution and set 75 mm above the capsules. 233 To determine the influence of sonication on the blue dextran re-234 235 lease, samples of about 250 capsules are sonicated in each test. The

236 total volume of blue dextran solution encapsulated in the capsules 237 is V_{in} . The capsules are placed in a volume V_{out} = 40 ml of solution C 238 in a 28 mm diameter tube.

239 2.4. Capsule compression

240 The capsule release induced by sonication is compared to that 241 induced by a more classical stimulation: mechanical stimulation 242 by compression. A computer-controlled traction/compression de-243 vice (Synergie 400, MTS Systems, France) is fitted with a 2 N force transducer (accuracy 10^{-4} N). A capsule filled with blue dextran is 244 placed on a lower plate within a transparent cup filled with solu-245 tion C (V_{out} = 10 ml). As shown in Fig. 1, it is compressed by a piston 246 that moves down at a constant speed. The piston velocity is set at 247 248 0.6 mm/min, which is low enough to eliminate inertia effects but large enough to avoid potential osmotic effects. At each time step, 249 250 the acquisition system records automatically the imposed displacement of the piston D(t) and the resultant force exerted on 251 252 the piston. The initial contact point between the piston and the capsule corresponds to D(0) = 0; it is determined with a precision 253 254 of ±20 μ m. We define the ratio $\delta(t) = D(t)/D_0$ as the compression 255 ratio. The buoyancy force acting on the piston is subtracted from

Fig. 1. Compression test of a capsule.

the force measured to determine the net force F acting on the capsule.

Different values of the maximum compression ratio are imposed: δ_{max} = 0.2, 0.4, 0.6, 0.8. For each test, 30 capsules are compressed one by one in the same external solution of volume V_{out} . The number of capsules tested has been chosen in order to reach values of blue dextran concentrations in the external solution that can be accurately measured with the spectrophotometer. The total volume of encapsulated blue dextran solution in the 30 capsules is V_{in} .

2.5. Measurement of capsule mechanical properties by compression

The capsule mechanical properties are obtained by compression following the method of Carin et al. [22]. Unloaded capsules are tested in the traction/compression device described above. The capsule mechanical properties are extracted from the experimental 270 curve of the reduced force F/r_0 versus $\delta(t)$ using the model of 271 Lardner and Pujara [27]. We consider three membrane constitutive 272 273 laws, the neo-Hookean, Skalak and Evans & Skalak laws, which correspond to different mechanical behaviors (see [22] for more 275 details). Once a membrane constitutive law is assumed, the surface area-dilatation modulus K is found for each value of δ by compar-276 ing the theoretical and measured forces. The constitutive law that corresponds to the rheological behavior of the alginate membrane is the one for which *K* remains constant with δ .

2.6. Measurement of the capsule release

The concentration of the external solution in blue dextran $C_r(t)$ is measured using a spectrophotometer (SPECORD S 300 UV VIS, Analytic Jena, Germany) at a 620 nm wavelength. The mass of blue dextran released $m_r(t)$ is then given by

$$m_r(t) = C_r(t) \times V_{out}.$$
 (1)

It is to be compared to the total mass initially encapsulated within the entire capsule sample m_0 . We will call mass release the ratio $m_r(t)/m_0$. In order to determine m_0 , the capsule samples are exposed to a high intensity sonication at the end of each release experiment. The sonication duration time is set in order to guarantee that all the capsules are ruptured. Measuring the concentration C_0 of the solution in blue dextran provides the value of the initially encapsulated mass:

 $m_0 = C_0 \times (V_{out} + V_{in}).$ (2)

3. Results and discussion

3.1. Capsule breakup induced by sonication

We have observed that sonication could lead to the capsule breakup depending on the time and power of sonication. We have therefore investigated the breakup threshold varying the conditions of sonication. For each test, the capsule final state (ruptured or unruptured) has been determined by naked eye. Fig. 2 is a log-log plot of the breakup threshold power P_s^{thr} as a function of the sonication time. Depending on the sonication power, breakup occurs in a matter of minutes (with a maximum of about 1 h). The plot shows that the lower the sonication power, the longer the sonication needs to be to reach the capsule breakup. It also indicates that the breakup threshold follows a negative power law. For a given sonication time, breakup is expected to occur for sonication powers larger than $P_s^{thr} = 483t_s^{-1.7}$ with P_s in Watts and t_s in minutes. This equation is an estimate of the breakup threshold.

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Fig. 2. Diagram of the capsule state at the end of sonication. The full symbols correspond to unruptured capsules and the empty symbols to ruptured capsules. The line is an estimate of the breakup threshold.

The capsule breakup can be attributed to a fatigue effect: it is a consequence of the vibrations generated on the capsule by the ultrasonic field [28,29]. The oscillation amplitude A_s is a function of the parameters of sonication:

$$A_{\rm s} = \frac{\sqrt{2P_{\rm s}/\rho cS}}{2\pi f},\tag{3}$$

323 where ρ is the density of the propagation medium, *c* is the speed of sound in the medium, *f* is the ultrasonic frequency and *S* is the sur-324 325 face area of the sonotrode. The vibration amplitude is only of a few 326 microns, but at such a high frequency, it is enough to generate fatigue for the thin elastic membrane of the capsule. The succession of 327 expansions and contractions of the capsule membrane eventually 328 329 leads to its rupture.

330 Breakup occurs when stresses are locally greater than the mate-331 rial yield stress. But since stresses and deformations are linked by 332 the local constitutive law, one can also consider that breakup 333 occurs when the total plastic deformation induced during the son-334 ication time t_s reaches a threshold value. For large oscillation 335 amplitudes and above the plastic threshold, the plastic deforma-336 tion after each cycle can be assumed to be proportional to the 337 amplitude A_s , itself proportional to $\sqrt{P_s}$ according to Eq. (3). Since 338 there are $t_s f$ cycles during the sonication time t_s , the total plastic 339 deformation is hence proportional to $t_s \sqrt{P_s}$. Such simple consider-340 ations predict that the breakup power is proportional to t_s^{-2} , which 341 is close to our observation. From the smaller exponent presently found, one may infer that the plastic deformation is actually pro-342 343 portional to the amplitude above a certain yield value.

344 3.2. Effect of sonication on the capsule mechanical properties

345 To analyze the influence of sonication on unruptured capsules, 346 we have measured the mechanical properties of sets of capsules 347 previously subjected to different conditions of sonication. We have first fixed the sonication power at 10.71 W and increased the son-348 ication time up to 9 min (breakup occurs at $t_s \sim 10$ min); then, set-349 350 ting the sonication time at 6 min, we have increased the sonication 351 power up to 17 W (breakup occurs at $P_s \sim 17.5$ W). The values of 352 the mechanical properties, obtained by compression, are compared 353 with measurements on non-sonicated capsules.

354 Prior to testing the capsules by compression, we have measured 355 their membrane thickness to see whether sonication had an influ-356 ence on it. The thickness is found not to change following sonica-357 tion, no matter what the parameters of the ultrasonic stimulation 358 are. We have therefore considered all the capsules to have a thick-359 ness $h_0 = 0.20 \pm 0.01$ mm.

360 When the capsules are subjected to compression, they are all 361 observed to recover their exact initial shape after 2 h without 362 any creeping effects. It proves that the effect of compression is

reversible and that the membrane has elastic properties. The small response time is mostly a consequence of the presence of the inner liquid.

To post-process the data curves obtained from the compression tests, we have used the inverse analysis of Carin et al. [22]. We have calculated the surface area-dilation modulus *K* as a function of the compression ratio for the three constitutive laws. We find that the neo-Hookean law is the only law for which it is, on the whole, independent of the compression ratio; large variations are otherwise obtained with the Skalak and Evans & Skalak laws. The capsules are therefore strain-softening. We find that the capsule alginate membrane follows the neo-Hookean law, whether the capsule is sonicated or not. It proves that the capsule recovers its elastic properties after sonication.

Fig. 3 shows the evolution of the apparent area-dilation modulus K/h_0 with the sonication parameters. For the neo-Hookean law, the area-dilation modulus is equal to the elastic or Young modulus: we will therefore call K/h_0 the apparent elastic modulus. No significant difference is found between non-sonicated and sonicated capsules, regardless of the sonication time and power. Below the breakup threshold, sonication therefore has a too small influence on the capsule mechanical properties for changes to be detected with the current measurement technique.

The alginate capsules are measured to have an apparent elastic modulus K/h_0 equal to 6 kPa on average. Even if no other study had previously characterized the mechanical properties of alginate capsules, the result can be compared with measurements obtained on alginate beads. Millimetric beads prepared with a 2.2% alginate solution were found to have a shear modulus just below 12 kPa, which corresponds to an elastic modulus of 36 kPa [30]. We therefore find coherent values of mechanical properties for the present capsules, which are fabricated with a 0.2% alginate solution: the difference in elastic modulus is a consequence of the very different concentrations in alginate and calcium ions used during both fabrication processes.

3.3. Effect of sonication on capsule release

3.3.1. Blue dextran release during sonication

The percentage of encapsulated mass that is released during the 400 time of sonication is studied for capsules filled with a large mole-401 cule of blue dextran. It is denoted $[m_r/m_0]$, where the brackets indi-402 cate the step increase. It is obtained for different conditions of 403 sonication by measuring the concentration in blue dextran of the 404 suspending solution before and after the sonication. Four sets of 405 capsules are subjected to a sonication power of 0.48 W for duration 406 times ranging from 1 min to 5 min. Another four sets of capsules 407 are sonicated for a fixed duration of 2 min under powers ranging 408 from 0.48 W to 3.06 W. Fig. 4 shows the evolution of the mass ratio 409 released during sonication for the different sets of capsules. We 410 have plotted the mass release as a function of the amplitude of 411 the ultrasonic pressure sine wave 412 413

$$p_s = \sqrt{2\rho c P_s/S} \tag{4}$$

instead of the power P_{s} . It shows that sonication leads to pressures of the order of a few bars. The mass release is found to be proportional to the sonication time (Fig. 4a) and ultrasonic pressure amplitude (Fig. 4b). The mass release is therefore also proportional to the ultrasonic \hat{w} ave amplitude A_s according to Eqs. (3) and (4). These relationships of proportionality indicate that, contrary to other types of stimuli, no threshold value exists for release to be induced by ultrasonic stimulation.

In order to provide a physical explanation for these results, we have explored alternative mechanisms to see which one accounts for a linear dependency between mass release and ultrasonic pres-

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Fig. 3. Apparent elastic modulus obtained with a neo-Hookean law for capsules subjected to different conditions of sonication. *P*_s = 0 W and *t*_s = 0 min correspond to the non-sonicated case.



Fig. 4. Step increase in release mass ratio (m_r/m_0) measured at the end of sonication as a function of the time of sonication t_s (a) and ultrasonic pressure p_s (b). $P_s = 0.48$ W corresponds to $p_s = 193.9$ kPa. The constants in the regression equations are respectively in minutes in (a) and in kPa in (b).

427 sure. One mechanism that may drive the internal liquid out is the 428 acoustic radiation force exerted by the acoustic waves on the cap-429 sules. If one assumes that the flow through the porous membrane 430 is governed by Darcy's law, one finds that the mass release is pro-431 portional to the radiation force. But since the radiation force is pro-432 portional to the acoustic power P_s [31], the mass release would then be proportional to p_s^2 (Eq. 4), which is not what is measured 433 experimentally. 434

435 Another possible mechanism is the acoustic streaming flow created by the ultrasonic wave. Analytical models have shown that for 436 a narrow sound beam, the streaming velocity U_s is proportional to 437 P_s in laminar flow conditions and to $\sqrt{P_s}$ at infinitely large Reynolds 438 numbers [32]. We have measured the velocity of the acoustic jet to 439 be about 70 cm/s in the case of a sonication power P_s = 3.08 W, 440 441 which is the maximum power tested in the mass release experi-442 ments. The capsules placed in the reservoir are entrained in the 443 large recirculating annular vortex around the jet, where they have a velocity of 5 cm/s. When they occasionally pass under the acous-444 tic jet, their velocity reaches 20 cm/s. The Reynolds number of the 445 446 flow around the millimetric capsules $Re = 2U_s r_0/v$ is therefore of the 447 order of $10^2 - 10^3$, meaning that the flow is laminar (the kinematic viscosity of the suspending fluid is $v = 10^{-6} \text{ m}^2/\text{s}$: the streaming 448 449 velocity U_s should be proportional to P_s and thus to p_s^2 . Such a scal-450 ing was found experimentally by Poindexter et al. [33]. The flow 451 around the capsules entrains the liquid by convection within the boundary layer that forms at the capsule surface. The exchange 452 flow induced by convection is determined by the boundary layer 453 thickness. In laminar flow conditions, the latter scales like $Re^{-1/2}$ 454 and hence like $U_s^{-1/2}$ [34]. Since the average mass flux across the 455 boundary layer is inversely proportional to its thickness, $[m_r/m_0]$ 456 scales like $U_s^{1/2}$ and is thus proportional to p_s . 457

458 In conclusion, convective effects induced by the acoustic 459 streaming flow account for the linear dependency between the

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mass release and the acoustic pressure. One can then also easily understand why the release is proportional to the sonication time: the longer the stimulation, the higher the release.

3.3.2. Comparison of the mass release induced by sonication and compression

For comparison, the mass release is measured for loaded capsules subjected to another type of mechanical stimulation: compression. The results are shown in Fig. 5a for different maximum compression ratios δ_{max} . The increase in released mass ratio is directly proportional to the maximum compression imposed on the capsule. The results indicate that, like for ultrasonic stimulation, release takes place even for the low values of compression: no threshold compression value needs to be reached for release to start. The results are also consistent with previous measurements on polyelectrolyte gels: a relationship of proportionality had been found in the special case of gels that are neutral in terms of ioncontent with the encapsulated liquid [35].

As the capsule is compressed, the pressure inside the capsule increases. The mass release is induced by the pressure gradient that builds up across the membrane. In order to compare the release induced by sonication and compression, we have plotted in Fig. 5b the mass release versus the pressure difference Δp corresponding to each value of δ_{max}

$$\Delta p = \frac{F_{max}}{S_{contact}}.$$
(5) 485

The force F_{max} is the average net force acting on the capsule at the compression ratio δ_{max} . The contact surface area between the piston and the compressed capsule $S_{contact}$ is obtained measuring the radius of the contact surface on the acquired images. The fact that the mass release is presently found to be proportional to $\sqrt{\Delta p}$ is a consequence of the mechanical behavior of the alginate

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Fig. 5. Step increase in mass ratio (m_r/m_0) released from capsules loaded with blue dextran under compression, as a function of the compression ratio maximum compression ratios δ_{max} (a) and corresponding pressure difference across the membrane Δp (b).

membrane. We have shown in Section 3.2 that the alginate capsule 492 membrane obeys to the neo-Hookean law. A numerical simulation 493 of the compression of a thick-membrane capsule with neo-Hook-494 495 ean properties has indeed proven that the force is roughly propor-496 tional to the square root of the compression ratio [21]. One must, however, keep in mind that the driven force (and hence mass re-497 498 lease) is lower if the capsule follows a less strain-softening or, even 499 more so, a strain-hardening constitutive law.

500 To reach a 40% release ratio, the capsule needs to be compressed 501 up to 67% of its initial height. Ultrasonically, the same value of 502 release is reached after a sonication time of only 2 min and 40 s 503 when using the lowest ultrasonic power tested $(P_{\rm s} = 0.48 \, \text{W})$. $p_s = 193.9$ kPa). Two advantages of the ultrasonic stimulation over 504 505 compression can be highlighted. The most obvious one is that sonication acts on the capsule at a distance and does not require a 506 direct contact with the capsule. Release can therefore be controlled 507 remotely. The second advantage is the very high accuracy that can 508 509 be achieved when controlling the mass release. Fig. 5a shows that it is difficult to release more than 50% of the encapsulated product 510 511 by compression. A full release is, however, possible by ultrasonic 512 stimulation.

513 3.3.3. Influence of sonication on the passive release

In order to investigate whether sonication has a permanent effect on the membrane porosity, we finally study the time-evolution of the passive release from loaded capsules. The passive release is measured on two sets of capsules just after their sonication $(P_s = 0.48 \text{ W}, t_s = 2 \text{ and } 5 \text{ min})$ and compared to the one measured on non-sonicated capsules (Fig. 6). Considering first the case without sonication, the curve that best fits the experimental points is



Fig. 6. Passive release of blue dextran for non-sonicated capsules (\blacktriangle) and capsules previously sonicated at 0.48 W for 2 min (\blacksquare) and 5 min (\bullet). *t* = 0 corresponds to the end of sonication.

$$\frac{m_{\rm r}(t)}{m_0} = 0.62 \left[1 - \exp\left(\frac{-t}{260}\right) \right],\tag{6}$$

in which the time constant is expressed in minutes. A nearly perfect 524 fit is found with the exponential law, the coefficient of determina-525 tion being $R^2 = 0.98$. The blue dextran release therefore tends expo-526 nentially toward the constant value of $(m_r/m_0)_{\infty} = 0.62$ with a 527 characteristic time constant of 260 min. A similar characteristic 528 time constant has been shown in the case of sodium polystyrene 529 sulfonate incorporated in alginate gels [36]. The good fit with the 530 exponential law proves that the release kinetics is diffusion-based: 531 Fick's first law, with its assumptions of homogeneous concentra-532 tions inside and outside the capsules and time-constant inner and 533 outer volumes, predicts a mass release 534

$$\frac{m_r(t)}{m_0} = \frac{V_{out}}{V_{in} + V_{out}} \left[1 - \exp\left(\frac{-\lambda A(V_{in} + V_{out})}{V_{in} V_{out}} t\right) \right],\tag{7}$$

where λ is the overall mass transfer coefficient and A the total exter-538 nal capsule surface area available for mass transfer [37]. Where the 539 blue dextran release diverges from Fick's law prediction is in the 540 asymptotic value reached at infinite times. The experimental value 541 of $(m_r/m_0)_{\infty}$ is much lower than the theoretical value $(V_{out}/m_0)_{\infty}$ 542 $(V_{in} + V_{out}) \sim 0.95$). With a molecular weight of 2000 kDa, the blue 543 dextran molecule has a Stoke's radius of \sim 27 nm, which is slightly 544 bigger than the typical alginate pore size [38]. The membrane there-545 fore limits the blue dextran diffusion, when the molecules travel 546 throughout. Previous studies have shown that, in such a case, the 547 encapsulated molecule is only partially released, the asymptotic 548 release ratio decreasing as the relative pore size is decreased [39]. 549

The passive release is then studied on sonicated capsules. Fig. 4a 550 indicates that a 2 min sonication at 0.48 W induces a released mass 551 ratio $[m_r/m_0] = 0.33$. Fig. 6 shows that the subsequent passive 552 release has a larger departure from the exponential curve 553 $(R^2 = 0.84)$. But on average, the release mass ratio still follows Eq. 554 (7). It tends toward the same constant $(m_r/m_0)_{\infty} = 0.62$ and the 555 characteristic time constant is almost identical (265 min). No sen-556 sible effect of sonication is therefore found on the passive release 557 properties of the capsule membrane. For a larger sonication time 558 (5 min), the initial released mass ratio is $[m_r/m_0] = 0.71$, which is 559 higher than $(m_r/m_0)_{\infty}$ (Fig. 4a). Fig. 6 shows that no subsequent 560 release or uptake takes place once the sonication is terminated. 561 All these results show that sonication does not have a remnant ef-562 fect on the membrane porosity, as the overall porosity returns back 563 to its original value after sonication. 564

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565 4. Conclusion

566 We have investigated the influence of sonication on capsules with a soft membrane made of hydrogel. No measurable effect 567 has been found on the capsule mechanical properties, as long as 568 the times and powers of sonication remain below a certain thresh-569 old. Above threshold, sonication leads to the capsule breakup be-570 571 cause of the fatigue of the membrane. When a substance is 572 encapsulated, sonication leads to an increase in the mass release. 573 The increase measured during sonication is found to be propor-574 tional to the duration time and pressure amplitude of the ultrasonic stimulation. This linear dependency can be explained by 575 576 acoustic streaming: the high-induced velocities enhance convection close to the capsule membrane and thus mass release. If one 577 subjects the capsule to a compression instead of an ultrasonic pres-578 sure, we also find a linear relationship of the increase in mass re-579 lease with the maximum compression ratio. But the comparison 580 581 proves the higher efficiency of sonication to reach large values of release ratios in short time instants. We have finally studied the 582 influence of sonication on the passive release to detect a possible 583 permanent effect on the membrane porosity. Sonication appears 584 585 to have no remnant effect, as the capsules recover their initial 586 properties on the whole.

587 The present study proves that ultrasonic stimulation could be 588 an effective method to remotely enhance the release from softmembrane liquid-filled alginate capsules or induce the capsule 589 590 breakup. By choosing the sonication time and power adequately, we have shown that it is possible to control the released dose of 591 the encapsulated molecule. A clinical use of sonication appears to 592 593 be feasible, as previous studies have shown that the range of ultra-594 sonic stimulation presently tested is harmless for cells [40,41]. 595 Sonication could therefore be a powerful external stimulation to trigger a local drug release from loaded polyelectrolyte capsules 596 in vivo. For this application, further studies would need to be con-597 ducted on smaller size capsules to prove that the results remain 598 valid, as vectorization mainly relies on micro- or nano-capsules. 599 600 The present results are, however, directly applicable to all the engi-601 neering processes that encapsulate fragile or volatile substances 602 and require their release at a precise instant of time. Sonication 603 should then be considered to remotely induced release, as it is a 604 very efficient technique to control a partial or total release within short time durations. 605

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